

# Defect mediated magnetic transitions in Fe and Mn doped MoS<sub>2</sub>

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Cite this: DOI: 10.1039/c8cp02882f

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We report single-phase syntheses of undoped 2H-MoS<sub>2</sub> as well as Mn and Fe doped MoS<sub>2</sub> by a facile hydrothermal route. The formation of the 2H-MoS<sub>2</sub> phase was confirmed by XRD and was corroborated with Raman spectra. The morphology of the doped and undoped MoS<sub>2</sub> nanostructures comprised sheets, as revealed by TEM and STEM images. The fine granular structure was observed by high resolution TEM micrographs. The STEM-EDS results show dopant concentrations of ~1 atom% corresponding to Mn and Fe in doped MoS<sub>2</sub>. The undoped MoS<sub>2</sub> revealed diamagnetic behavior at room temperature and paramagnetic behavior in the range (100 to 300 K). The Mn-MoS<sub>2</sub> sample displayed ferromagnetism below 20 K with a coercive field of ~50 O<sub>e</sub>. Such a sample may be utilized for magnetic switching purposes at low temperatures. The onset of the antiferromagnetic interaction was observed below 145 K in Fe-MoS<sub>2</sub> samples. They have been understood in terms of long-range magnetic interactions amongst the dipole moments mediated *via* surface defects as well as the interaction between the dipoles and the surface charges. The findings are corroborated with the help of EPR studies.

 Received 6th May 2018,  
Accepted 17th May 2018

DOI: 10.1039/c8cp02882f

rsc.li/pccp

## 1. Introduction

Transition metal dichalcogenides (TMDCs) have received considerable attention from researchers owing to their applications in the areas of optoelectronics, electronics, photo catalysts, magnetic switching and storage.<sup>1–9</sup> Molybdenum disulfide (MoS<sub>2</sub>) in particular with a direct band gap is a suitable candidate for optoelectronic and photo catalytic applications.<sup>10–12</sup> The ease of doping and heterostructure formation with tunable and desired properties make it an attractive alternative to graphene.<sup>13</sup> 2H-MoS<sub>2</sub> in its bulk form is semiconducting in nature and doping with transition metals such as Fe, Ni, Co, and Mn makes it a promising candidate for dilute magnetic semiconductors (DMS).<sup>4,9,14–18</sup> Nanostructural manipulations *via* grain and grain boundary engineering might open an avenue to tailor the magnetic properties of transition metal doped MoS<sub>2</sub>.<sup>19</sup> In this study, we have synthesized MoS<sub>2</sub>, Fe and Mn doped MoS<sub>2</sub> by a hydrothermal method.

Unlike previous studies where MoS<sub>2</sub> is reported to be ferromagnetic in nature,<sup>3,5,15,20–22</sup> we could synthesize single-phase MoS<sub>2</sub> powders with a diamagnetic character. It may be noted that ferromagnetic characteristics in MoS<sub>2</sub> have been

observed for thin films deposited over substrates. In this investigation, we also demonstrate the effect of Fe and Mn dopants in MoS<sub>2</sub> on their magnetic behavior. It is well known that in their atomic/ionic form Fe and Mn display long-range ferromagnetic and anti-ferromagnetic ordering, respectively, owing to the Heisenberg exchange interaction. However, we found reversal of their magnetic behaviors in doped samples of MoS<sub>2</sub>. This is perhaps the most interesting part of our investigation. The undoped MoS<sub>2</sub> revealed diamagnetic behavior at room temperature and paramagnetic behavior in the range (100 to 300 K). The Mn-MoS<sub>2</sub> sample displayed ferromagnetism below 20 K with a coercive field of ~50 O<sub>e</sub>. The onset of antiferromagnetic interaction was observed below 145 K in Fe-MoS<sub>2</sub> samples. This has been understood in terms of long-range magnetic interactions amongst the dipole moments mediated *via* surface defects as well as the interaction between the dipoles and the surface charges. It will be shown that magnetic behavior in this investigation has arisen thanks to doping, as we are dealing with a single-phase material with structural nano domains of similar lengths.

## 2. Synthesis methods

The MoS<sub>2</sub> nanostructures were synthesized by a hydrothermal method. In this, 1 millimole of ammonium heptamolybdate (AHM) was mixed in 30 mL of de-ionized water and stirred

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1 magnetically to get a clear solution. Into this solution 28  
 2 millimoles of thiourea were added and stirred until complete  
 3 dissolution. The above solution was transferred into a 50 mL  
 4 Teflon-lined hydrothermal autoclave reactor and placed in a  
 5 Muffle furnace at 220 °C for 18 h. The reaction product was  
 6 cleaned by centrifugation and ultrasonication cycles using DI  
 7 water and finally with ethanol. The product was vacuum dried  
 8 at 60 °C to obtain a black-colored powder. FeCl<sub>3</sub>·3H<sub>2</sub>O and  
 9 MnSO<sub>4</sub> precursors from Alfa Aesar were utilized for the syn-  
 10 thesis of Fe and Mn doped MoS<sub>2</sub>, respectively. The doping of ~3  
 11 atom per cent of these transition metals required the addition  
 12 of 0.21 millimoles of dopant precursors into AHM and thiourea  
 13 solution. The resulting solutions were stirred magnetically to  
 14 achieve homogeneity. The remaining procedures were the same  
 15 as those for the synthesis of MoS<sub>2</sub>.

### 3. Instrumentation

20 The prepared undoped, Mn and Fe doped MoS<sub>2</sub> powders were  
 21 investigated at room temperature using a PANalytical X. Pert  
 22 Pro-X-ray diffractometer (XRD) with Ni-filtered Cu K<sub>α</sub> radiation  
 23 of wavelength 1.54 Å. A Renishaw inVia RM2000 Raman spec-  
 24 trometer was used to collect the Raman spectra for undoped,  
 25 Mn, and Fe doped MoS<sub>2</sub> using a 785 nm excitation wavelength,  
 26 a 50× objective, and 10 seconds exposure time; the power used  
 27 was 1 per cent. Transmission electron microscopy (TEM) in-  
 28 vestigations and scanning TEM-energy dispersive X-ray spectro-  
 29 scopy (STEM-EDS) characterization were carried out with FEI,  
 30 Tecnai G<sup>2</sup>T20 and Talos F200X microscopes operating at 200  
 31 kV. Room-temperature electron paramagnetic resonance (EPR)  
 32 measurements were performed on a Bruker A200X-band CW  
 33 spectrometer using 9.83 GHz microwave frequency, 8 mW  
 34 microwave power and 10 G modulation amplitude. The EPR  
 35 parameters were evaluated by simulating the experimental  
 36 spectra using EASYSPIN.<sup>23</sup> Magnetic hysteresis and suscepti-  
 37 bility were recorded using Quantum Design SQUID-VSM. The  
 38 magnetic field was zeroed before every measurement from an  
 39 initial +10 Oe in an oscillating mode at room temperature. This  
 40 was done in order to minimize the effects of a trapped magnetic  
 41 field in the SQUID.

42 The phase and the nature of the layered structure were  
 43 determined by XRD and Raman spectroscopy, respectively.  
 44 The morphologies and lattice structures were investigated with  
 45 TEM and electron diffraction patterns. The STEM-EDS was  
 46 carried out for micro-chemical analysis of Fe and Mn doped  
 47 MoS<sub>2</sub> samples. The magnetic response was studied in terms of  
 48 magnetization by varying the temperature from 300 K to 5 K.  
 49 The data was collected in the Zero Field Cooling (ZFC) and Field  
 50 Cooled Cooling (FC) modes of measurement.

## 4. Results and discussion

### 4.1. Structural characterization

55 Fig. 1 shows the XRD pattern of undoped, Mn doped, and Fe  
 56 doped MoS<sub>2</sub> nanostructures with the standard XRD pattern of

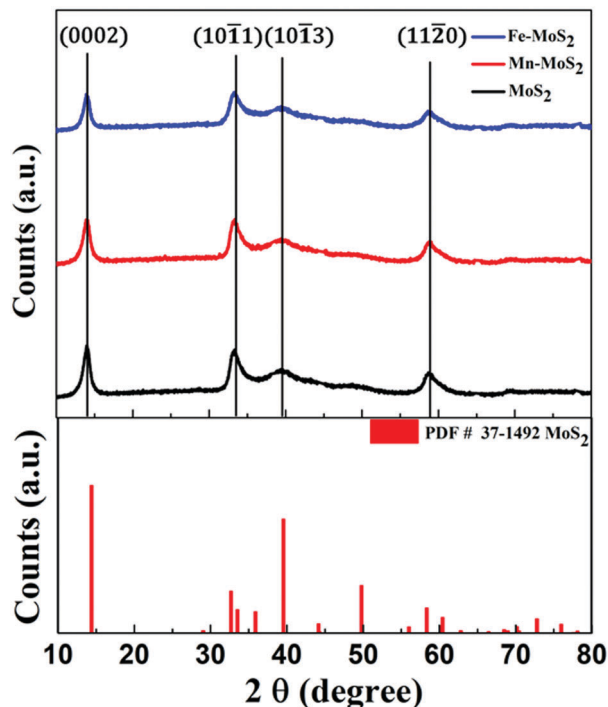


Fig. 1 Powder X-ray diffraction patterns of MoS<sub>2</sub>, Mn doped MoS<sub>2</sub> and Fe doped MoS<sub>2</sub> nanostructures. The bottom figure shows the standard XRD patterns of 2H-MoS<sub>2</sub> (PDF #37-1492).

MoS<sub>2</sub> (bottom figure). The reflections (0002), (10 $\bar{1}$ 1), (10 $\bar{1}$ 3) and (11 $\bar{2}$ 0) corresponding to 2H-MoS<sub>2</sub> (PDF #37-1492) were observed. The absence of additional peaks in the XRD pattern of Mn and Fe doped MoS<sub>2</sub> rule out the possibilities of formation of their sulfide phases. Moreover, the Mn and Fe doped MoS<sub>2</sub> patterns do not show a significant change in the disposition of the peaks compared to that of undoped MoS<sub>2</sub>. This may be attributed to the dilute addition of Mn and Fe dopants in the MoS<sub>2</sub>. It may be noted that under the present synthesis conditions (with variable pH and S<sup>2-</sup> concentrations) sulfide formation of the dopants is more favorable compared to their unreacted soluble salts. This in turn rules out the possibility of removal of dopants ions during the subsequent washing protocol. Therefore, dopants are incorporated into the MoS<sub>2</sub> lattice and a single-phase Mn and Fe doped MoS<sub>2</sub> has formed. The Scherrer formula was used to measure the crystallite size of the undoped and doped MoS<sub>2</sub> samples. The crystallite sizes corresponding to the (10 $\bar{1}$ 1), (10 $\bar{1}$ 3) and (11 $\bar{2}$ 0) reflections of the undoped, Mn, and Fe doped MoS<sub>2</sub> samples are estimated to be ~8.0 nm, 5 nm, 3 nm, and 3 nm, respectively. Since MoS<sub>2</sub> is a layered material, as confirmed by our TEM results (to be presented later), and (0002) is the axis normal to the layer, which suggests that the average thickness of the as-synthesized nanostructures is ~7 monolayers. Fig. 2 displays the Raman spectra of undoped, Mn doped and Fe doped MoS<sub>2</sub> nanostructures. The characteristics peaks at the ~384 cm<sup>-1</sup> and ~408 cm<sup>-1</sup> positions correspond to the E<sub>2g</sub> and A<sub>1g</sub> vibration modes of the phonon, suggesting that the undoped synthesized MoS<sub>2</sub> nanostructures is 2H-MoS<sub>2</sub>.<sup>24</sup> Doping with Mn and Fe in MoS<sub>2</sub> did not

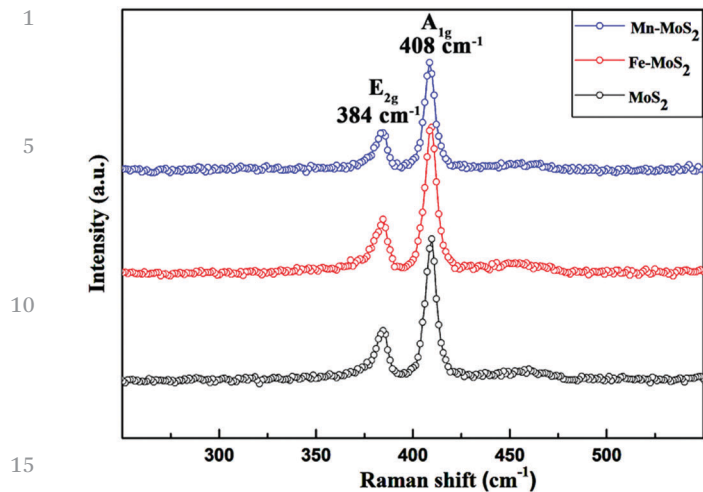


Fig. 2 Raman spectra of the undoped, Mn doped, and Fe doped  $\text{MoS}_2$  single 2H- $\text{MoS}_2$  phase. The characteristic peaks at positions at  $\sim 384 \text{ cm}^{-1}$  and  $\sim 408 \text{ cm}^{-1}$  corresponding to the  $E_{2g}$  and  $A_{1g}$  modes of phonon vibrations are shown.

reveal any additional peak in the Raman spectra conforming once again to the dilute addition of dopants into  $\text{MoS}_2$ .<sup>9</sup> From the XRD and Raman results, we conclude that both undoped and doped  $\text{MoS}_2$  nanostructures display a single 2H- $\text{MoS}_2$  phase.

Fig. 3 shows a high resolution TEM (HRTEM) micrograph of Mn doped  $\text{MoS}_2$ , with insets of the Fast Fourier Transform (FFT) in the bottom right as well as a magnified view of the selected region in the top left. The presence of a thin layer and

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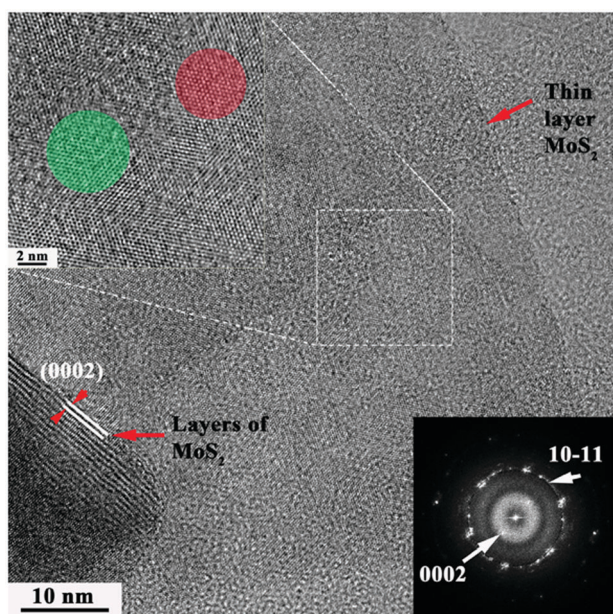


Fig. 3 HRTEM micrograph with FFT as inset (bottom right) displays spots corresponding to the (0002) and (10 $\bar{1}$ 1) planes of 2H- $\text{MoS}_2$ . Arrows (top right and bottom left) show thin layers of 2H- $\text{MoS}_2$ . The magnified HRTEM image (inset top left) of the selected region of this (white square) depicts the crystalline domains. Two such domains are highlighted using false color. The (0002) plane of 2H- $\text{MoS}_2$  is also shown in the figure.

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several layers of Mn doped  $\text{MoS}_2$  is shown with arrows. The FFT revealing reflections corresponding to (10 $\bar{1}$ 1) of 2H- $\text{MoS}_2$  is depicted in the micrograph. The magnified image from the selected region shows nanocrystalline domains with an average size of  $\sim 1.3 \pm 0.3 \text{ nm}$  (averaged over at least 35 regions chosen statistically). Two such domains are depicted by putting false colors (red and green) on them.

Fig. 4 displays the STEM-HAADF-EDS elemental maps of Mo, S, and Mn. The HAADF image (Fig. 4a) shows the sheet-like morphology of the Mn- $\text{MoS}_2$  nanostructures. Maps display the uniform distribution of Mo, S, and Mn. Further, the absence of clusters in the elemental map of Mn rules out the possibility of Mn-sulfide phases. Therefore, Mn has been distributed uniformly in 2H- $\text{MoS}_2$  and a single phase has formed. The elemental compositions of Mn, Mo, and S were found to be  $\sim 1$  atom%,  $\sim 32$  atom%, and  $\sim 65$  atom%, respectively.

The bright field TEM image and corresponding selected area diffraction pattern of Fe doped  $\text{MoS}_2$  are shown in Fig. 5. A flower-like morphology was observed and is comprised of buckled sheets of  $\text{MoS}_2$ . The diffraction rings are indexed as (10 $\bar{1}$ 1), (10 $\bar{1}$ 3) and (0008) reflections of 2H- $\text{MoS}_2$  phases (PDF #37-1492). The appearance of diffuse rings indicates the presence of fine crystalline nano domains in the sample. Fig. 6 displays the STEM-HAADF-EDS elemental maps of Mo, S, and Fe in Fe- $\text{MoS}_2$  in Fig. 6c, d and e, respectively. The distributions of the elements are uniform in the sheets, conforming to the single phase. The HRTEM micrographs of the Mn and Fe doped  $\text{MoS}_2$  are shown in Fig. 7. The FFTs (given in the inset) are comprised of streaks in the spots, indicating that the nanocrystalline structural domains are misoriented from each other but the angle of misorientation appears to be very small. Hence, nanocrystalline domains are observed in both Mn and Fe doped  $\text{MoS}_2$ . Unlike all previous reports, we are dealing with the dilute addition of antiferromagnetic (Mn) and

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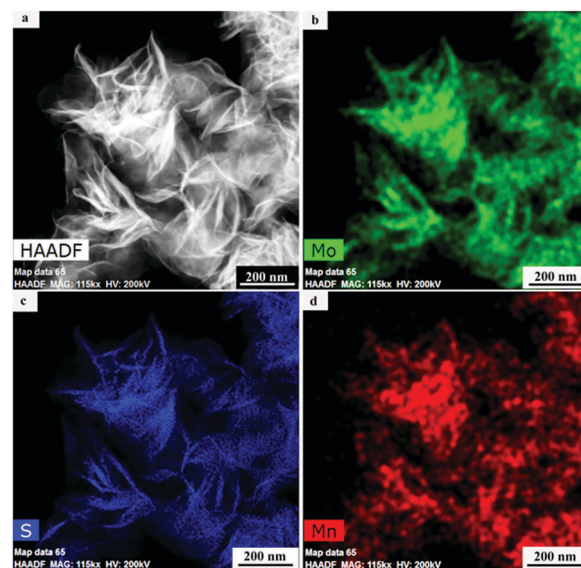


Fig. 4 STEM-HAADF-EDS elemental mapping of Mn doped  $\text{MoS}_2$ . High angle annular dark field (HAADF) image (a) and elemental maps corresponding to Mo (b), S (c), and Mn (d), respectively.

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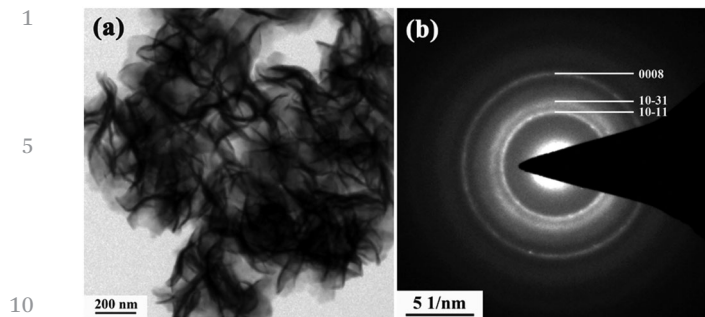


Fig. 5 Bright field TEM micrograph of Fe doped MoS<sub>2</sub> (a) and corresponding selected area diffraction pattern (b).

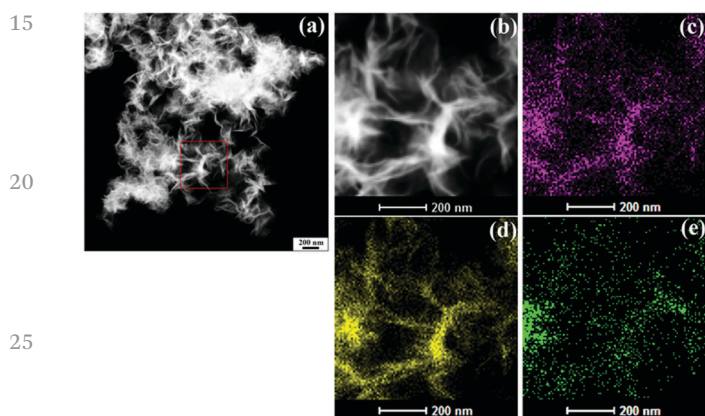


Fig. 6 Scanning TEM-HAADF image of Fe doped MoS<sub>2</sub> (a), HAADF image of the selected area (b), elemental maps corresponding to Mo (c), S (d), and Fe (e), respectively.

4 QS

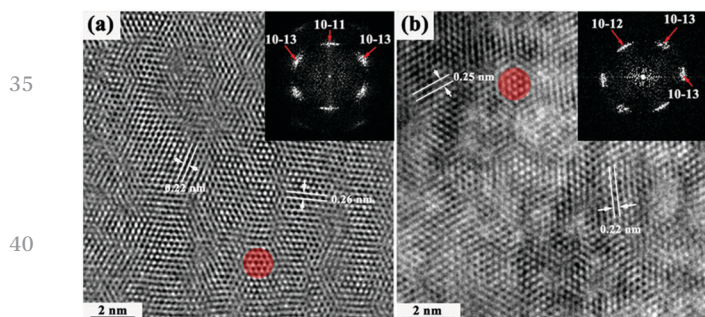


Fig. 7 HRTEM micrographs of Mn doped MoS<sub>2</sub> (a) and Fe doped MoS<sub>2</sub> (b) with their FFTs as insets. The nano domains are shown by applying false colors in the micrographs.

ferromagnetic (Fe) elements as well as a single phase. Hence, the formation of sulfide phases is ruled out *a priori*. In view of this, we conclude that all the dopants have gone into the planar hexagonal net of MoS<sub>2</sub>. This was substantiated by our XRD, Raman spectra, HRTEM and STEM-EDS results.

#### 4.2. Magnetic characterization

The magnetic response of MoS<sub>2</sub> powder used in the present study shows diamagnetic behavior (*cf.* Fig. 8a). Unlike previous

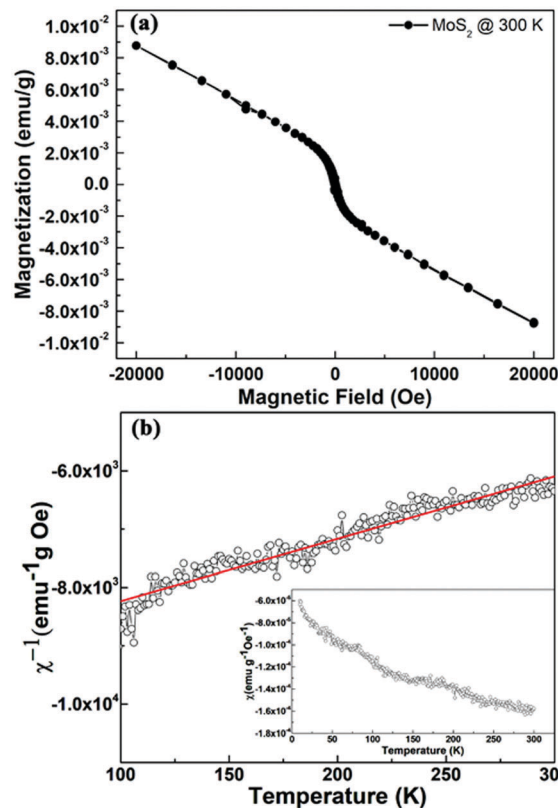


Fig. 8 (a) Magnetization versus magnetic field loop for MoS<sub>2</sub> at 300 K revealing the diamagnetic character of the sample, (b) inverse of magnetic susceptibility ( $\chi^{-1}$ ) versus temperature of MoS<sub>2</sub> powder where the data fits linearly with the Curie law. The inset shows the magnetic susceptibility ( $\chi$ ) versus temperature plot of MoS<sub>2</sub> powder measured at 100 Oe.

reports where a ferromagnetic hysteresis loop is evident on top of the diamagnetic background in MoS<sub>2</sub> films, this powder presented pure diamagnetic behavior without any long-range ordering. This could be attributed to the negligible ferromagnetic signal arising from the edges or vacancies in the bulk limit of MoS<sub>2</sub>. The diamagnetic susceptibility of MoS<sub>2</sub> is also examined by varying the temperature and is plotted in Fig. 8(b). Unlike a pure diamagnetic material where temperature-independent diamagnetic susceptibility is observed, we find temperature-dependent behavior, which is shown in Fig. 8(b). By plotting the inverse of the susceptibility ( $\chi$ ) as a function of temperature, linear Curie-Weiss law behavior is observed between 100 K and 300 K. This could be ascribed to the presence of a few paramagnetic centers in MoS<sub>2</sub>. As for the fitting parameters, the Curie constant  $C$  is found to be  $9.6 \times 10^{-4}$  emu K g<sup>-1</sup> Oe<sup>-1</sup> and the Weiss temperature  $T_c$  is found to be  $\sim 885$  K. The presence of a high and positive  $T_c$  in MoS<sub>2</sub> is in agreement with previous reports where magnetization is observed due to diamagnetic, paramagnetic and ferromagnetic interactions arising from different types of defects.<sup>25</sup> The result is in conformity with the weak EPR lines of the MoS<sub>2</sub> sample exhibiting uniaxial  $g$ -anisotropy with  $g_{\parallel} = 2.010$  and  $g_{\perp} = 2.002$  (see Fig. 9). Similar  $g$ -values have been reported previously for MoS<sub>2</sub> samples and attributed to sulfur (S)-related defects at the

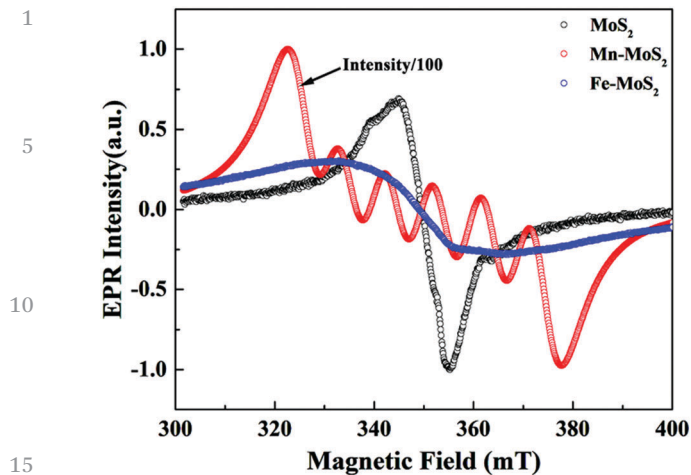


Fig. 9 EPR spectra of undoped, Mn, and Fe doped MoS<sub>2</sub> samples. The undoped MoS<sub>2</sub> displays uniaxial anisotropy in  $g$ , Mn-MoS<sub>2</sub> depicts an enhanced signal accompanied by hyperfine splitting, and Fe-MoS<sub>2</sub> shows isotropy in  $g$  with reduced signal intensity. The arrow shows the EPR signal corresponding to Mn-MoS<sub>2</sub> which is scaled by 100 to fit into the intensity range of the undoped and Fe-MoS<sub>2</sub> samples.

poor symmetry sites/edges of nanocrystalline MoS<sub>2</sub>. The other high field weak EPR lines could be due to Mo-related local defects arising either from the dangling bonds or dislocations.<sup>26</sup> It may be pointed out that the TEM micrographs do show a sheet-like morphology in this investigation. However, magnetic measurements from powder samples mask its individual contribution.

The magnetization of Mn-MoS<sub>2</sub> is plotted as a function of magnetic field at different temperatures in Fig. 10(a). It can be noted from the figure that at 300 K, the sample is diamagnetic whereas it becomes paramagnetic at 50 K and by further lowering the temperature to 20 K, it behaves like a ferromagnet with a small coercive field of  $\sim 50$  Oe. In order to understand the thermal ordering in the sample, the temperature-dependent susceptibility of Mn-MoS<sub>2</sub> and its inverse are given in Fig. 10(b) and (c). Three dominant phase transitions can be noted from the figure. These refer to the temperature ranges (i) 23–300 K, where the magnetic susceptibility is fairly small and constant, (ii) 40–230 K, where the susceptibility increases by lowering the temperature (Curie–Weiss law) and (iii) below 35 K, where a sharp increase in  $\chi$  is observed for both ZFC and FC modes. Further, to elaborate the observed transitions, the temperature range (70–230 K) is fitted with the Curie–Weiss law (see Fig. 10(c)). The observed Curie constant  $C$  is  $\sim 6.43 \times 10^{-5}$  emu K  $g^{-1}$  Oe<sup>-1</sup> and  $T_c$  is  $\sim 55$  K. The  $T_c$  value suggests that there exists a weak ferromagnetic phase along with the paramagnetic background in Mn-MoS<sub>2</sub>. As the free Mn ions are paramagnetic by nature, the presence of weak ferromagnetism can be warranted due to exchange interaction between free electrons on Mn ions and surface charges on MoS<sub>2</sub> arising due to S-vacancy.<sup>9</sup> These charges align around Mn-ions and thus give rise to the ferromagnetic coupling. Interestingly, by further lowering the temperature, the  $\chi$  value bifurcates for ZFC and FC curves at 30 K. Due to the presence of uncompensated surface

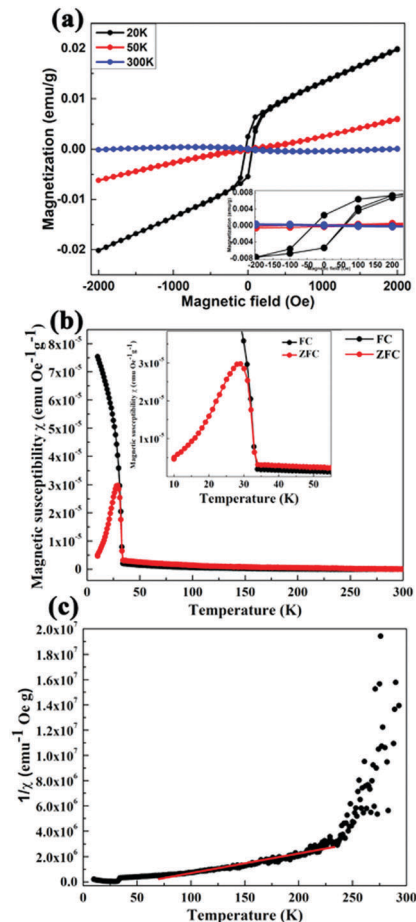


Fig. 10 (a) Magnetization versus magnetic field loop for Mn-MoS<sub>2</sub> revealing the diamagnetic character of the sample at 300 K, paramagnetism at 50 K and ferromagnetism at 20 K. The inset of the figure shows the enlarged hysteresis loop of the sample in the weak field regime. (b) Magnetic susceptibility ( $\chi$ ) versus temperature (between 5 K and 300 K) of Mn doped MoS<sub>2</sub> powder in both ZFC and FC cooled mode at 100 Oe and (c)  $1/\chi$  versus  $T$  plot in the temperature range 5–300 K where the Curie–Weiss law is fitted in the range 70–230 K.

charges at low temperatures, bifurcation occurs, revealing the onset of another long-range interaction amongst Mn-ions and the surface charges. As reported previously, for nano domain sizes smaller than the magnetic exchange length, the local anisotropies are smoothed by the exchange interactions such that their effects are negligible on total magnetization. For such cases, in general, small coercive fields are observed.<sup>27,28</sup> The observed small coercivity in our case conforms to this, making Mn-MoS<sub>2</sub> a suitable soft magnet for magnetization reversal and memory storage applications.

In order to elucidate the room-temperature paramagnetic properties, the EPR spectrum of Mn-MoS<sub>2</sub> is also examined and is shown in Fig. 9. Like previous reports, we also observed the presence of six characteristic hyperfine lines of Mn<sup>2+</sup> ( $I = 5/2$ ) with axially anisotropic  $g$ -values 2.011 and 2.003 and hyperfine interaction constants 283 and 81 MHz.<sup>9,29</sup> The broad peak corresponding to S-vacancy is also observed in the background. Interestingly, the EPR signal is nearly two orders stronger in

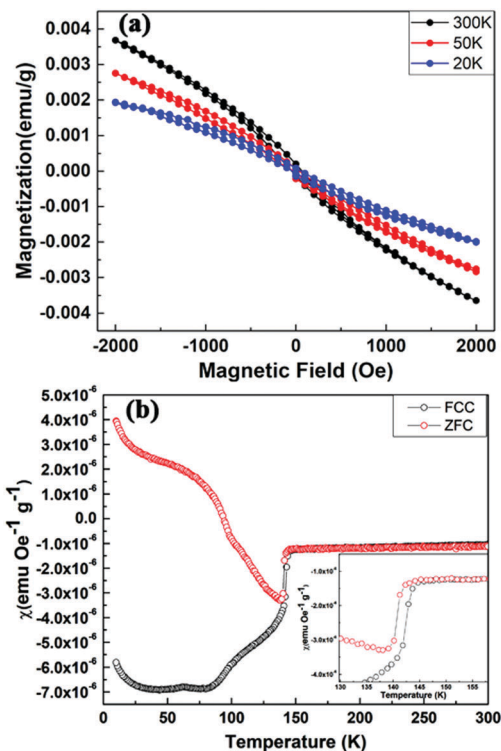


Fig. 11 (a) Magnetization versus magnetic field loop for Fe doped MoS<sub>2</sub> revealing the diamagnetic character of the sample at 20 K, 50 K and 300 K, (b) Magnetic susceptibility ( $\chi$ ) versus temperature (between 5 K and 300 K) of Fe-MoS<sub>2</sub> powder in both ZFC and FC mode at 100 Oe. The inset shows magnetic susceptibility ( $\chi$ ) versus temperature (between 120 K and 160 K) plot of Fe-MoS<sub>2</sub> powder, revealing the transition temperature around 145 K.

Mn-MoS<sub>2</sub> than that in MoS<sub>2</sub>, suggesting that the surface charge densities on MoS<sub>2</sub> have been increased by the Mn substitution. These results corroborate the existence of a ferromagnetic interaction between the defect-related surface charges and the dopants at low temperatures. These results are in agreement with those of the XRD and Raman spectra, suggesting the substitution of Mn ions over the MoS<sub>2</sub> surface in a single phase. STEM-EDS results also substantiate these findings.

The room-temperature EPR spectrum of Fe-MoS<sub>2</sub> reveals the presence of a broad peak with an isotropic  $g$ -value  $g = 2.015$  (*cf.* Fig. 9). The intensity of the EPR signal is nearly half of that of MoS<sub>2</sub>, suggesting the randomization of the surface charges on MoS<sub>2</sub> and their local interactions with the dopants. The magnetization of the Fe-MoS<sub>2</sub> sample is also plotted as a function of magnetic field at different temperatures in Fig. 11(a). The sample shows diamagnetic behavior at 20 K, 50 K and 300 K with variable slopes and small coercive fields ( $\sim 20$  Oe). The observed behavior cannot be classified as purely diamagnetic nor weakly ferromagnetic, as both characters are evident from the hysteresis loops. In addition, lower saturation magnetization for Fe-MoS<sub>2</sub> is observed than that of Mn-MoS<sub>2</sub>, which confirms the randomization of surface charges. There are relatively finer structural domains in Fe-MoS<sub>2</sub> ( $\sim 1.1$  nm) than Mn-MoS<sub>2</sub> ( $\sim 1.3$  nm) at the given temperature.<sup>30,31</sup> The weak

magnetization in Fe-MoS<sub>2</sub> can be attributed to the preferred substitution of Fe ions on the S-edge over the Mo-edge of the MoS<sub>2</sub> sheets and the dominant diamagnetic background of the MoS<sub>2</sub> matrix.<sup>8</sup>

The temperature-dependent magnetic susceptibility of Fe-MoS<sub>2</sub> is plotted in both ZFC and FC modes and is shown in Fig. 11(b). The sample remains mainly diamagnetic in the temperature range 150–300 K. However, by further lowering the temperature, a sharp decrease in  $\chi$  and its thermal hysteresis are noted at 145 K, revealing the onset of antiferromagnetic interactions amongst Fe ions and the local matrix. Below 138 K, the  $\chi$  value bifurcates for the ZFC and FC modes. We attribute these thermally irreversible differences in the magnetization to the coexistence of antiferromagnetic and ferromagnetic interactions in the sample.<sup>32,33</sup> It is worth-mentioning that the effects of higher magnetic fields on the bifurcation magnitude and temperature need to be explored for a better understanding of the observed behavior. Previously, the small and negative ZFC moments under weak magnetic fields have been attributed as artefacts due to trapped magnetic fields.<sup>34</sup> In our case, the field was zeroed before every measurement from an initial +10 Oe in an oscillating mode at room temperature (as mentioned in section 3). According to Kumar and Yousuf,<sup>34</sup> this may introduce negatively trapped fields and consequently a negative ZFC magnetization could be observed. However, in our case, we have rather observed an increase in the ZFC magnetization than that of FC mode. We therefore rule out the possibility of any artefacts being present in the sample due to trapped magnetic fields.<sup>34,35</sup> The  $\chi$  in ZFC mode is greater than that in the FC mode for Fe-MoS<sub>2</sub>. This reveals that cooling of the sample under a magnetic field (100 Oe) is aligning the moments in the opposite direction to the applied magnetic field and thus a lower value of magnetization is observed for the FC mode than that of the ZFC mode. Further, we are dealing with a single-phase material with structural nano domains of  $\sim 1.0$  nm. This permits us to infer that the dilute addition of Mn and Fe is responsible for the observed magnetic responses. The extremely fine structural nano domains and substitution of dopants in the 2H-MoS<sub>2</sub> lattice create defects as well as surface charges. The complex interplay of defects and surface charges is leading to long-range magnetic interactions amongst the dipoles.

## 5. Conclusions

We have synthesized undoped and, Mn and Fe doped MoS<sub>2</sub> nanostructures by a hydrothermal method. The phase of undoped MoS<sub>2</sub> was found to be 2H-MoS<sub>2</sub>, as revealed by XRD and Raman spectra. The absence of sulfide phases in Mn and Fe doped samples were confirmed by XRD, Raman spectra, and STEM-EDS results. The EPR result shows the presence of Mn<sup>2+</sup> ions in the Mn-doped MoS<sub>2</sub> sample. The undoped MoS<sub>2</sub> showed diamagnetic behavior at room temperature convoluted with a feeble ferromagnetic response. It shows a paramagnetic nature in the temperature window (100 K to 300 K) with  $T_c \sim 884$  K.



1 The presence of a high and positive  $T_c$  in MoS<sub>2</sub> could be  
 attributed to the diamagnetic, paramagnetic and ferromagnetic  
 interactions arising from different types of defects. The Mn-  
 MoS<sub>2</sub> revealed ferromagnetic behavior (below 20 K) with a small  
 5 coercive field of  $\sim 50$  O<sub>e</sub>. The Fe-MoS<sub>2</sub> sample exhibited the  
 onset of an antiferromagnetic interaction amongst Fe ions and  
 the local matrix below 145 K. Thus, we have demonstrated a  
 reversal in the magnetic response of Mn and Fe through their  
 dilute addition in MoS<sub>2</sub> *vis-à-vis* their elemental magnetic  
 10 behaviors. Thus, research on the dilute addition of different  
 types of magnetic ions in MoS<sub>2</sub> holds promise for potential  
 applications in magnetic switching at low temperature.

## 15 Conflicts of interest

There are no conflicts to declare.

## 20 Acknowledgements

Two of us (BM and PC) acknowledge the financial support  
 under DST-Inspire faculty and student schemes.

## 25 Notes and references

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